



# Crystal structures and microwave dielectric properties of novel low-permittivity $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$ ceramics

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## ABSTRACT

$\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2\text{--}1.0$ ) microwave dielectric ceramics were synthesized using the solid-state method at the temperature from 1087 °C to 1150 °C for 3 h. All compositions showed a single phase for  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  when  $x$  increased from 0.2 to 1.0, and a phase transition from monoclinic to triclinic structure occurred between 0.8 and 1.0. The relative permittivity of  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ceramics decreased from 6.57 to 6.12 when the  $\text{Ba}^{2+}$  ions substituted by  $\text{Sr}^{2+}$  ions. However, the quality factor initially decreased from 34,735 GHz ( $x = 0.2$ ) to 28,986 GHz ( $x = 0.4$ ) and subsequently increased monotonously. The variation in the temperature coefficient of resonant frequency showed an opposite, slightly changing trend compared with that of the quality factor. A novel single-phase  $\text{SrZnSi}_3\text{O}_8$ , which possesses good microwave dielectric properties of  $\epsilon_r = 6.12$ ,  $Q \times f = 78,064$  GHz, and  $\tau_f = -33.2$  ppm/°C, was obtained for the first time at the sintering temperature of 1150 °C.

## 1. Introduction

With the expanding operating frequency ranges of microwave wireless communication, wireless communication is on the verge of entering its fifth generation. High frequency is required for ultrahigh-speed local area networks, electronic toll systems, and car anti-collision systems based on the intelligent transport systems [1,2]. Low-permittivity ( $\epsilon_r < 15$ ) microwave dielectric ceramics can be used as high-frequency substrates, dielectric antennae, high-accuracy capacitors, and millimeter-wave components, such as resonators and filters [3]. Thus, high-performance microwave dielectric ceramics with low permittivity is expected to attract considerable attention. For millimeter-wave wireless communication, microwave dielectric ceramics must present a low dielectric constant ( $\epsilon_r$ ) to reduce the transmission attenuation and the cross-coupling effect, a high-quality factor ( $Q \times f$ ) to achieve excellent frequency selectivity, and a near-zero temperature coefficient of resonant frequency ( $\tau_f$ ) to ensure the stability of the transmitted frequency [4,5].

Silicates exhibit several interesting characteristics, such as ferroelectricity in  $\text{Bi}_2\text{SiO}_5$  and  $\text{BaZnSiO}_4$  [6,7], unique thermal expansion coefficient in  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_2\text{O}_7$  [8], and microwave dielectric properties

in  $\text{Mg}_2\text{Al}_4\text{Si}_5\text{O}_{18}$  [9]. These interesting features are related to the materials' complex crystal structure consisting  $[\text{SiO}_4]$  tetrahedrons and other polyhedral. In general, silicates present a low permittivity value because of the Si–O bond in  $[\text{SiO}_4]$  tetrahedron, which contains 45% ionic bond and 55% covalent bond; the covalent bond reduces the relative permittivity because of the decrease of rattling effect [1]. In addition, numerous kinds of silicates, such as  $\text{Zn}_2\text{SiO}_4$ ,  $(\text{Sr}_{1-x}\text{A}_x)_2(\text{Zn}_{1-x}\text{B}_x)\text{Si}_2\text{O}_7$  ( $A = \text{Ca}, \text{Ba}$  and  $B = \text{Co}, \text{Mg}, \text{Mn}, \text{Ni}$ ) and  $\text{CaAl}_2\text{Si}_2\text{O}_8$  [10–12], have been explored. All of these silicates present a low permittivity value and a high quality factor, rendering the materials as good candidates for millimeter-wave devices.

In recent years, there have been several works focusing on feldspar-based microwave dielectric ceramics, which comprise plagioclase-feldspar-type and alkali-feldspar-type structures. The chemical formulas are  $A[\text{B}_2\text{C}_2]\text{O}_8$  ( $A = \text{Ca}, \text{Sr}, \text{Ba}$ ;  $B = \text{Al}, \text{Ga}$ ;  $C = \text{Si}, \text{Ge}$ ) and  $A[\text{BC}_3]\text{O}_8$  ( $A = \text{K}, \text{Na}$ ;  $B = \text{Al}, \text{Ga}$ ;  $C = \text{Si}, \text{Ge}$ ), respectively [12–16]. The feldspar crystal structure is composed of a three-dimensional framework of  $[\text{Al}/\text{GaO}_4]$  or  $[\text{Si}/\text{GeO}_4]$  tetrahedral. The charge balance due to the substitution of  $\text{Al}^{3+}/\text{Ga}^{3+}$  for  $\text{Si}^{4+}/\text{Ge}^{4+}$  is maintained by the addition of an interstitial alkali or alkali-earth ion [13]. Except for the most common feldspar, that is, zinc-feldspar, is rarely noticed. In 1970,

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Segnit et al. [17] synthesized the  $\text{BaZnSi}_3\text{O}_8$  ceramics by the solid-state reaction method. Heuer et al. [18] reported the crystal structure of single-crystal  $\text{CaZnSi}_3\text{O}_8$  prepared by hydrothermal method in 1998. Fehr et al. [19] called  $\text{CaZnSi}_3\text{O}_8$  as zinc-feldspar and predicted the existence of  $\text{SrZnSi}_3\text{O}_8$ . In our previous work,  $\text{BaZnSi}_3\text{O}_8$  ceramics was synthesized and its microwave dielectric properties were reported [3]. Thus so far,  $\text{SrZnSi}_3\text{O}_8$  has not yet been synthesized, and its microwave dielectric properties have not been reported previously. In this work,  $\text{Sr}^{2+}$  ions were used as substitute for  $\text{Ba}^{2+}$  ions, and  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-0.8$ ) solid solution was prepared through the solid-state reaction method. Besides, the phase composition, microstructure, and microwave dielectric properties of  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-1.0$ ) ceramics were investigated.

## 2. Experimental procedure

The  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-1.0$ ) ceramics were prepared by conventional solid-state method using reagent grade  $\text{BaCO}_3$  (99.8%),  $\text{SrCO}_3$  (99.8%),  $\text{ZnO}$  (99.5%),  $\text{SiO}_2$  (99.5%) powder as raw materials. According to desired stoichiometry, the raw materials were weighed to ball milled in a polyethylene jar for 12 h using  $\text{ZrO}_2$  balls with deionized water. After drying at  $85^\circ\text{C}$ , the mixtures were calcined in air at  $1000^\circ\text{C}$  for 3 h with a heat rate of  $5^\circ\text{C}/\text{min}$ . And then the powders were uniaxially pressed into samples with dimensions of 12 mm in diameter and approximately 6 mm in height under a pressure of 150 MPa. The samples were sintered in the temperature range of  $1087^\circ\text{C}$ – $1150^\circ\text{C}$  for 3 h at a heating rate of  $5^\circ\text{C}/\text{min}$ , and then they were naturally cooled in the furnace after being cooled to  $1000^\circ\text{C}$  at a rate of  $1^\circ\text{C}/\text{min}$ .

The apparent density of the sintered samples was measured by Archimedes' method. The relative density  $\rho_{\text{rel}}$  was obtained by as formula follows [20]:

$$\rho_{\text{rel}} = \frac{\rho_{\text{app}}}{\rho_{\text{the}}} \quad (1)$$

where  $\rho_{\text{app}}$  and  $\rho_{\text{the}}$  were the apparent density and theoretical densities, respectively.

The X-ray diffraction (XRD) data were obtained using the XRD-7000 X-ray diffractometer (Shimadzu, Kyoto, Japan) with  $\text{CuK}\alpha$  radiation. The microstructure and elemental mapping of the  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  samples were observed by scanning electron microscope (Sirion 200, Netherlands) with an energy-dispersive X-ray spectrometer (JSM-6490LV, Japan). Grain size distributions were obtained using the Image J software. The values of  $\epsilon_r$  and the unloaded  $Q \times f$  value were measured within 12–14 GHz in the  $\text{TE}_{011}$  mode by Hakki and Coleman method [21] using a network analyzer (Agilent E8362B, Agilent Technologies, USA) and parallel silver boards. The  $\tau_f$  value in the temperature range of  $30$ – $80^\circ\text{C}$  was calculated by Eq. (2):

$$\tau_f = \frac{1}{f(T_0)} \frac{[f(T_1) - f(T_0)]}{T_1 - T_0} \quad (2)$$

where  $f(T_1)$  and  $f(T_0)$  represented the resonant frequency at  $T_1$  ( $80^\circ\text{C}$ ) and  $T_0$  ( $30^\circ\text{C}$ ), respectively.

## 3. Results and discussion

Fig. 1 (a) shows the XRD patterns of  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-1.0$ ) ceramics sintered at their optimum temperature. The diffraction peaks that corresponded to  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-0.8$ ) are indexed to  $\text{BaZnSi}_3\text{O}_8$  (PDF#23-0841), indicating that  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  forms a solid solution with monoclinic structure (P21/m space group). However, for  $x = 1.0$ ,  $\text{SrZnSi}_3\text{O}_8$  is not assigned a PDF card. The lattice parameters of  $\text{BaZnSi}_3\text{O}_8$  and  $\text{SrZnSi}_3\text{O}_8$  were extracted from XRD data using the least-squares method. All the peaks are indexed accordingly, and the crystal structure information is given in the supplemental file.

Fig. 1 (b) presents the enlargement of the main diffraction peaks

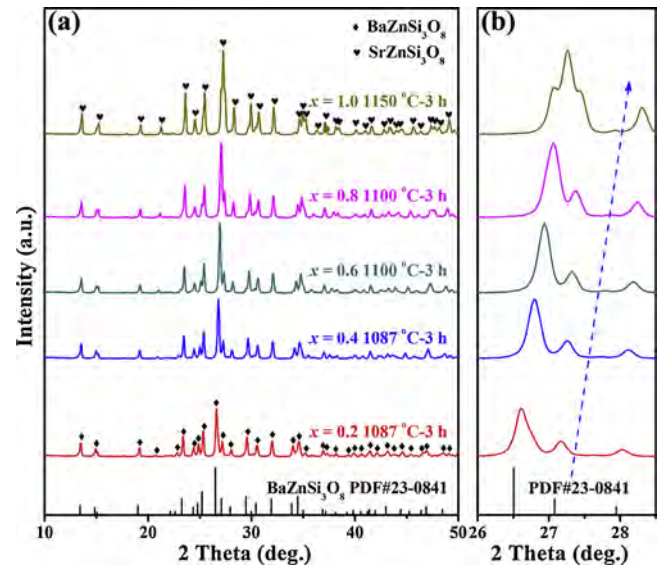


Fig. 1. (a) XRD powder diffraction patterns of  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-1.0$ ) ceramics sintered at  $1087^\circ\text{C}$ – $1150^\circ\text{C}$  for 3 h. (b) Enlargement of the main diffraction peaks between  $26.0^\circ$  and  $28.5^\circ$ .

between  $26.0^\circ$  and  $28.5^\circ$ . Resulting from the substitution of  $\text{Sr}^{2+}$  ions for  $\text{Ba}^{2+}$  ions, the peaks shifts to high angles. This phenomenon is expected because the ionic radius of  $\text{Sr}^{2+}$  is  $1.31 \text{ \AA}$  (CN = 9), which is smaller than that of  $\text{Ba}^{2+}$  ions ( $1.47 \text{ \AA}$ ; CN = 9) [22]. Moreover, the two main diffraction peaks at approximately  $26.5^\circ$  approaches each other when  $x$  increases from 0.2 to 1.0. However, another diffraction peak appears on the left of the main diffraction peak when  $x = 1.0$ . This phenomenon is caused by the phase transition from the P21/m space group ( $x = 0.2-0.8$ ) to the P-1 space group ( $x = 1.0$ ), as verified from the indexed results in the supplemental file. Meanwhile, the intensity of the diffraction peaks gradually increase as  $x$  increased, corresponding to the increase in crystallinity in the matrix ceramics [9].

Fig. 2 (a1)–(e1) present the microstructure of thermally etched  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-1.0$ ) ceramics sintered at  $1087^\circ\text{C}$  to  $1150^\circ\text{C}$  for 3 h. All of the compositions exhibit a dense microstructure, and small pores can be observed on the surface. Fig. 2 (a2)–(e2) and (a3)–(e3) show the backscattered electron (BSE) images and mapping of  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-1.0$ ) ceramics sintered at  $1087^\circ\text{C}$  to  $1150^\circ\text{C}$  for 3 h. No distinct grain can be observed (Fig. 2 (a2)–(e2)), and the elements of  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-1.0$ ) ceramics (Fig. 2 (a3)–(e3)) are evenly distributed on the surface, indicating that single-phase  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-1.0$ ) is formed. The average grain sizes estimated from Fig. 2 (a4)–(e4) are approximately 0.616, 0.723, 0.796, 0.874 and  $1.098 \mu\text{m}$ , corresponding to  $x = 0.2-1.0$ . The increase in average grain size corresponds to the enhancement of the main diffraction peaks in Fig. 1, indicating an increase in crystallinity [9].

Fig. 3 shows the apparent densities and relative densities of  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-1.0$ ) ceramics sintered at  $1087^\circ\text{C}$  to  $1150^\circ\text{C}$  for 3 h. The apparent densities of  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-1.0$ ) ceramics decreases monotonously from  $3.55 \text{ g/cm}^3$  to  $3.30 \text{ g/cm}^3$  because the atomic weight of Sr is smaller than that of Ba. The relative densities of all compositions are higher than 95%, and the effect of pores on the microwave dielectric properties is small [23].

Fig. 4 (a) shows the microwave dielectric properties of  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-1.0$ ) ceramics sintered at  $1087^\circ\text{C}$  to  $1150^\circ\text{C}$  for 3 h. The influence of porosity on  $\epsilon_r$  is eliminated by the following equation [24]:

$$\epsilon_{r-\text{corr}} = \epsilon_r (1 + 1.5P) \quad (3)$$

where  $P$  is the porosity and can be calculated by the formula as follows:

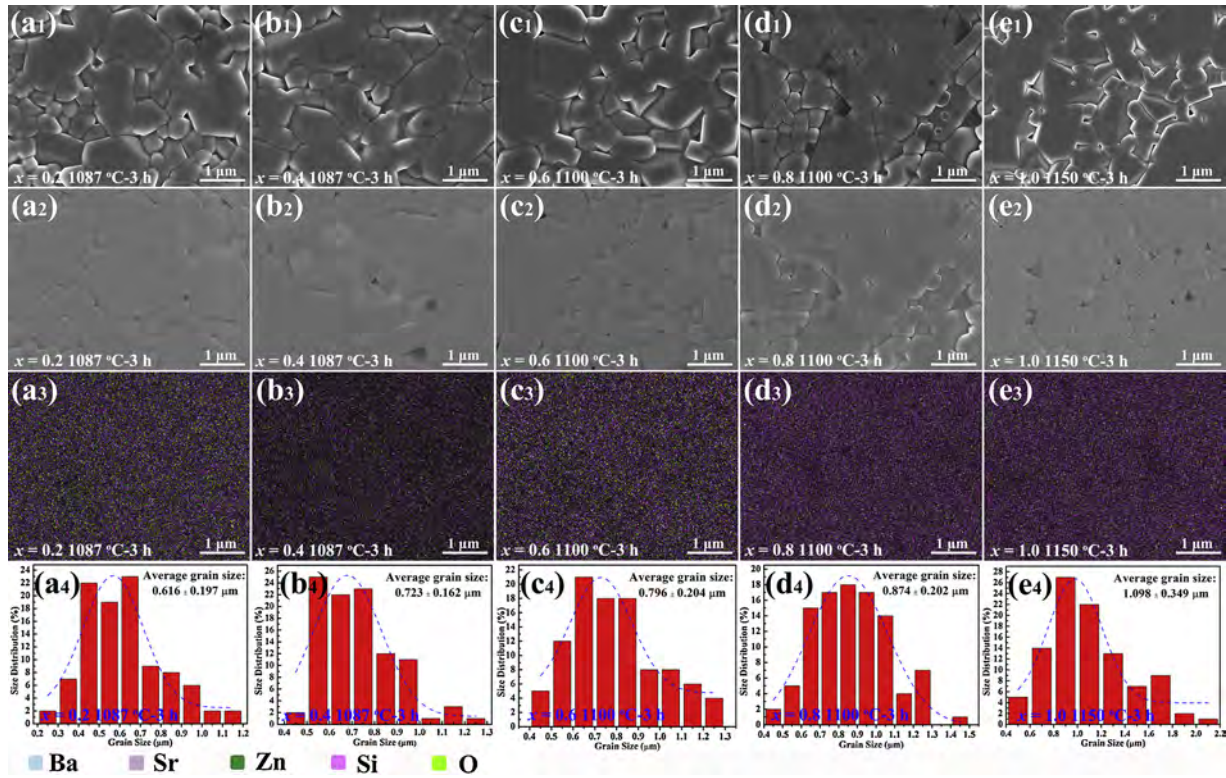


Fig. 2. (a1)-(e1) SEM images, (a2)-(e2) BSE images, (a3)-(e3) Mapping and (a4)-(e4) Grain size distribution of thermally etched  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-1.0$ ) ceramics sintered at  $1087^\circ\text{C}$ – $1150^\circ\text{C}$  for 3 h.

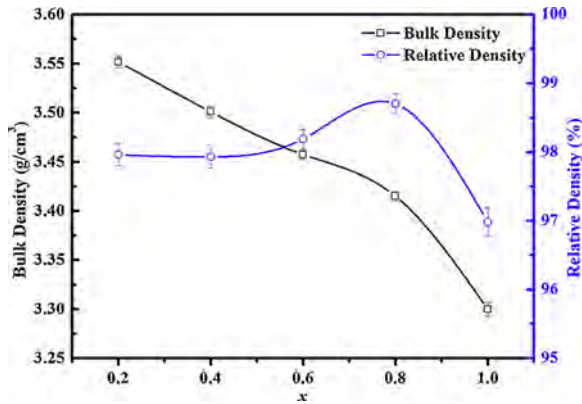


Fig. 3. Apparent densities and relative densities of  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-1.0$ ) ceramics sintered at  $1087^\circ\text{C}$ – $1150^\circ\text{C}$  for 3 h.

$$P = 1 - \rho_{\text{rel}} \quad (4)$$

The total ionic polarizability ( $\alpha_D^T$ ) of  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-1.0$ ) can be calculated by the additive rule [25,26]:

$$\alpha_D^T = (1-x)\alpha(\text{Ba}^{2+}) + x\alpha(\text{Sr}^{2+}) + \alpha(\text{Zn}^{2+}) + 3\alpha(\text{Si}^{4+}) + 8\alpha(\text{O}^{2-}) \quad (5)$$

where  $\alpha(\text{Ba}^{2+})$ ,  $\alpha(\text{Sr}^{2+})$ ,  $\alpha(\text{Zn}^{2+})$ ,  $\alpha(\text{Si}^{4+})$  and  $\alpha(\text{O}^{2-})$  are the ionic polarizabilities of  $\text{Ba}^{2+}$ ,  $\text{Sr}^{2+}$ ,  $\text{Zn}^{2+}$ ,  $\text{Si}^{4+}$ , and  $\text{O}^{2-}$ , respectively.

The relative permittivity ( $\epsilon_r$ ), porosity-corrected relative permittivity ( $\epsilon_{r-\text{corr}}$ ) and ionic polarizability ( $\alpha_{\text{theo}}$ ) of  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ceramics decreases monotonously when  $x$  increases from 0.2 to 1.0, thereby indicating that the ionic polarizability of  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  plays a dominant role in controlling the relative permittivity [26,27]. The  $Q \times f$  value initially decreases from 34,736 GHz for  $x = 0.2$  to 28,986 GHz for  $x = 0.4$ , and subsequently increases to 78,064 GHz for  $x = 1.0$  (Fig. 4 (b)). Variation of  $Q \times f$  values of  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  is similar to that of

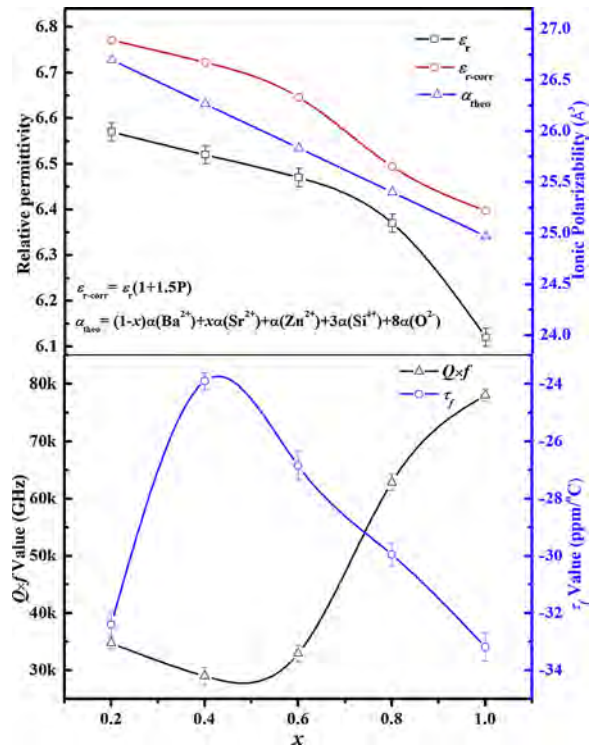


Fig. 4. (a) Relative permittivity ( $\epsilon_r$ ) and dielectric polarizability ( $\alpha_{\text{theo}}$ ) (b)  $Q \times f$  value and  $\tau_f$  value of  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2-1.0$ ) ceramics sintered at  $1087^\circ\text{C}$ – $1150^\circ\text{C}$  for 3 h.

( $\text{Sr}_{1-x}\text{A}_x$ )<sub>2</sub>ZnSi<sub>2</sub>O<sub>7</sub> ( $A = \text{Ca}, \text{Ba}$ ) reported by M.T. Sebastian in 2010 [11]. They attributed this phenomenon to the difference in radius of the A-site cations, which leads to synergistic effects, such as the orderly

arrangement of the cations.  $\text{Ba}_{1-x}\text{Sr}_x\text{Al}_2\text{Si}_2\text{O}_8$  and  $\text{Ba}_{1-x}\text{Ca}_x\text{Al}_2\text{Si}_2\text{O}_8$  solid solutions belong to the typical feldspar-type crystal structure, which have a similar crystal structure to  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$ . And the variation of  $Q \times f$  for  $\text{Ba}_{1-x}\text{Sr}_x\text{Al}_2\text{Si}_2\text{O}_8$ ,  $\text{Ba}_{1-x}\text{Ca}_x\text{Al}_2\text{Si}_2\text{O}_8$  and  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  solid solutions have the same trend [12,15]. Feldspar crystal structure is composed of a three-dimensional framework of  $[\text{Al}/\text{GaO}_4]$  or  $[\text{Si}/\text{GeO}_4]$  tetrahedral. It is easy to form a short-ordered arrangement Si and Al and long-ordered arrangement can be obtained after a long period of annealing. However, this Si:Al ordering cannot be detected by XRD due to the small difference in the X-ray scattering efficiencies of Al and Si. In  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  solid solution, the three-dimensional framework is composed of  $[\text{ZnO}_4]$  and  $[\text{SiO}_4]$  tetrahedral. It is more likely to form a short-ordered arrangement between  $\text{Zn}^{2+}$  and  $\text{Si}^{4+}$  due to their large difference of ionic radius and valence state. Krzmann attributes the variation of  $Q \times f$  for  $\text{Ba}_{1-x}\text{Sr}_x\text{Al}_2\text{Si}_2\text{O}_8$  and  $\text{Ba}_{1-x}\text{Ca}_x\text{Al}_2\text{Si}_2\text{O}_8$  solid solutions to the change of cation ordering, which may also apply to  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  solid solution. The variation trend of  $\tau_f$  values and  $Q \times f$  values is the opposite. The  $\tau_f$  values vary slightly between  $-23.9$  ppm/ $^\circ\text{C}$  and  $-33.2$  ppm/ $^\circ\text{C}$  in the entire range of  $x$ . The most excellent microwave dielectric properties (namely,  $\epsilon_r = 6.12$ ,  $Q \times f = 78,064$  GHz, and  $\tau_f = -33.2$  ppm/ $^\circ\text{C}$ ) are obtained when  $x = 1.0$  ( $\text{SrZnSi}_3\text{O}_8$ ).

#### 4. Conclusions

Low-permittivity  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  ( $x = 0.2$ – $1.0$ ) low-permittivity microwave dielectric ceramics have been synthesized using the conventional solid-state method. XRD patterns, BSE images and mapping show that  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  solid solutions with monoclinic structure (P21/m space group) are formed when  $x$  increases from 0.2 to 0.8. Importantly, a novel single-phase  $\text{SrZnSi}_3\text{O}_8$  ceramic with triclinic structure (P-1 space group) is obtained for the first time. The relative permittivity of  $\text{Ba}_{1-x}\text{Sr}_x\text{ZnSi}_3\text{O}_8$  is dominated by ionic polarizability, and the  $\epsilon_r$  decreases linearly from 6.57 to 6.12 when the  $\text{Ba}^{2+}$  ions are substituted by  $\text{Sr}^{2+}$  ions. However, the quality factor initially decreases from 34,735 GHz ( $x = 0.2$ ) to 28,986 GHz ( $x = 0.4$ ) and subsequently increases monotonously after a certain point. The  $\tau_f$  values varies slightly between  $-23.9$  ppm/ $^\circ\text{C}$  and  $-33.2$  ppm/ $^\circ\text{C}$  in the total  $x$  range. The most excellent microwave dielectric properties (namely,  $\epsilon_r = 6.12$ ,  $Q \times f = 78,064$  GHz, and  $\tau_f = -33.2$  ppm/ $^\circ\text{C}$ ) are obtained when  $x = 1.0$  ( $\text{SrZnSi}_3\text{O}_8$ ).

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